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# **Rare Earth and High Field Strength Element Partitioning Between Iron-Rich Clinopyroxenes and Felsic Liquids**

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#### **Abstract**

Rare earth elements are commonly assumed to substitute only for Ca in clinopyroxene because of the similarity of ionic radii for  $REE^{3+}$  and  $Ca^{2+}$  in 8-fold coordination. The assumption is valid for Mg-rich clinopyroxenes for which observed mineral/melt partition coefficients are readily predicted by the lattice strain model for substitution onto a single site (e.g. Wood and Blundy, 1997). We show that natural Fe-rich pyroxenes in both silica-undersaturated and –oversaturated magmatic systems deviate from this behavior. Salites (Mg# 48 to 59) in phonolites from Tenerife, ferrohedenbergites (Mg# 14.2 to 16.2) from the rhyolitic Bandelier Tuff, and ferroaugites (Mg# 9.6 to 32) from the rhyolitic Rattlesnake Tuff have higher heavy REE contents than predicted by single-site substitution. The ionic radius of  $Fe^{2+}$  in 6fold coordination is substantially greater than that of  $Mg^{2+}$ , hence we propose that, in Fe-rich clinopyroxenes, heavy REE are significantly partitioned between 8-fold Ca sites and 6-fold Fe-Mg sites such that Yb and Lu exist dominantly in 6-fold coordination. We also outline a REE-based method of identifying pyroxene-melt pairs in systems with multiple liquid and crystal populations, based upon the assumption that LREE and MREE reside exclusively in 8-fold coordination in pyroxene.

Contrary to expectations, interpolation of mineral/melt partition coefficient data for heavy REE does not predict the behavior of Y. We speculate that mass fractionation effects play a role in mineral/melt lithophile trace element partitioning that is detectable among pairs of isovalent elements with near-identical radii, such as Y-Ho, Zr-Hf, and Nb-Ta.

#### **Introduction**

Clinopyroxene is one of the most abundant minerals in igneous rocks and plays a major role in the generation and subsequent differentiation of magma. Trace element partitioning between clinopyroxene and silicate liquid has therefore been the subject of intense study, especially since the successful application of lattice strain theory (Brice, 1975; Blundy and Wood, 1994) to crystal-melt element partitioning in igneous systems (Gaetani and Grove, 1995; Wood and Blundy, 1997, 2001, 2002; Blundy et al., 1998; Lundstrom et al., 1998; Hill et al., 2000; Pertermann and Hirschmann, 2002). Most of these studies focus on Mg-rich clinopyroxene and its role in influencing lithophile trace element abundances in basaltic magma. Consequently, the partitioning of rare earth elements into Mg-rich clinopyroxene is well characterized: light REE are incompatible, but compatibility increases with atomic number, typically flattening out between Ho and Lu (Fig. 1). In contrast, Fe-rich clinopyroxenes from many strongly differentiated systems show somewhat different behavior; the value of the partition coefficient *D*  $[=(\text{concentration in solid phase})/(\text{concentration in liquid}), C_s/C_l]$  rises with atomic number from the light to the middle REE, flattens out or declines slightly between Gd and Ho-Er, and distinctly rises to higher values at Yb and Lu (Fig. 1). Here, we describe enhanced heavy REE contents in natural Fe-rich clinopyroxene from phonolites and rhyolites. In Mg-rich clinopyroxene, REE are assumed to reside exclusively in 8-fold coordination on the M2 structural site. We show that the distinctive partitioning behavior of REE into Fe-rich pyroxenes is most simply explained by partial heavy REE occupancy of 6-

fold coordinated sites with radii near those of  $Fe^{2+}$  and  $Mn^{2+}$ ; these may be located on either the M2 or M1 site (Cameron and Papike, 1980). We then describe a method for identifying equilibrium pyroxene-glass pairs on the basis of the complementary assumption that light and middle REE reside exclusively in 8-fold coordination. We also demonstrate significant Y/Ho, Zr/Hf and Nb/Ta fractionation between clinopyroxene and melt, and speculate that factors other than ionic radius and charge have a detectable effect on mineralmelt partitioning behavior of lithophile element cations.

### **Samples**

Pyroxenes and glasses in this study come from three volcanic systems: Na-salites from the phonolitic Fasnia Member of the Diego Hernández Formation, Tenerife, Spain; ferrohedenbergites from the highsilica rhyolitic Bandelier Tuff, New Mexico, USA; and ferroaugites to ferrohedenbergites from the rhyolitic Rattlesnake Tuff, Oregon, USA.

Information about each system is summarized in Table 1. The three magmas are all close to minimum melt compositions with  $(Na + K)/A \approx 1$  and near-identical degrees of melt polymerization as indicated by the parameter NBO/T, but vary strongly in alkali content and silica saturation.

# *Fasnia Member, Tenerife*

The 309  $\pm$  6 ka Fasnia Member (FM) is one of the major units of the phonolitic, caldera-related Diego Hernández Formation of the Las Cañadas Upper Group on Tenerife (Martí et al., 1994; Edgar et al., 2007). Las Cañadas phonolites have whole-rock compositions that lie close to the 1 kbar water-saturated phonolitic minimum in the *ne-ks-Q* system (Hamilton & MacKenzie, 1965), oxygen fugacities that lie close to the NNO buffer (Wolff & Storey, 1983) and temperatures typically in the range 790 – 850°C (Wolff & Storey, 1983; Bryan, 2006; Andújar et al., 2008). In addition to the phenocryst phases listed in Table 1, the FM contains trace amounts of apatite and pyrrhotite as inclusions in other minerals. The FM is the product the explosive eruption of  $\sim$ 13 km<sup>3</sup> of phonolite magma with a large range in trace element abundances (Edgar et al., 2007), due to mixing of at least two distinct end member phonolites and a small volume of mafic magma that resulted in a dominant volume of hybrid phonolite (Olin, 2007). Despite variable trace element contents (Table 2), fresh FM glasses considered here exhibit very little variation in major elements, with approximately 15 wt.% total alkalies and  $(Na + K)/A \approx 1$ .

Sodian salites analyzed in this study have near-constant major element compositions (Fig. 2) and trace element contents that are consistent with having crystallized from the hybrid FM magma (Olin, 2007). Individual pyroxene grains are euhedral and do not display detectable zoning in either major or trace elements.

# *Bandelier Tuff*

The  $\sim$ 450 km<sup>3</sup> Otowi Member of the Bandelier Tuff (BT) erupted from the Valles caldera, New Mexico, at 1.61 Ma. The BT is a high-silica rhyolite with a composition that closely corresponds to the water-saturated granite minimum between 1 and 2 kbar (Tuttle  $\&$  Bowen, 1958). The phenocryst assemblage contains quartz, magnetite and fayalite and  $fO<sub>2</sub>$  is therefore presumed to lie close to the OMF buffer. Trace amounts of zircon and allanite occur in addition to the phenocrysts listed in Table 1, and chevkinite and britholite are present as inclusions in clinopyroxene, fayalite, and magnetite (Min et al., 2006). The tuff displays strong vertical zonation in incompatible trace elements (Spell et al., 1990; Wolff et al., 1999); euhedral ferrohedenbergites and glasses analyzed in this study come from the most differentiated, chemically homogeneous, first-erupted portion of the magma body.

# *Rattlesnake Tuff*

The 7.05 Ma Rattlesnake Tuff (RT) represents the explosive eruption of  $\sim$ 280 km<sup>3</sup> of multi-component rhyolitic magma (Streck & Grunder, 1997, 2007). The RT contains five distinct high silica rhyolite compositions, A, B, C, D, and E (Streck  $\&$  Grunder 1997), related primarily by fractional crystallization where type A rhyolite is the most differentiated and type E is the least differentiated. Glasses and

pyroxenes from RT rhyolites type C, D, and E were analyzed in this study, and due to their similarities in trace element contents, respectively (Streck & Grunder, 1997), are treated as one population. In addition to the phases listed in Table 1, trace amounts of biotite occur in rhyolite types C - E, plus the accessory minerals zircon, chevkinite, apatite and pyrrhotite. The RT pyroxenes span a major-element compositional range from Mg# = 9.6 associated with type E rhyolite with  $\sim$ 74 wt% SiO, to pyroxenes with values as high as Mg# = 31.7 in type C rhyolite with ~77 wt%  $SiO<sub>2</sub>$ .

### **Analytical Methods**

All data were generated in the Washington State University GeoAnalytical Laboratory using laser ablation and solution ICP-MS analysis of trace element concentrations and electron microprobe for major element analysis, except for BT hedenbergite major element microprobe data from Kuentz (1986). Trace element concentrations in FM, BT and RT pyroxenes and FM and BT glasses were measured by ablating shallow (<5 µm deep) troughs with a New Wave 213 nm laser coupled to an Element2 high-resolution mass spectrometer. For pyroxenes, ablation troughs were 12 to 20  $\mu$ m wide by 0.5 to 1 mm long. Grains were analyzed *in situ* where encountered in polished sections (for FM and BT samples) or as mineral separates mounted in epoxy (for FM, BT and RT samples). Trace element concentrations in pumiceous glasses from the FM and BT were determined using LA-ICP-MS *in situ* on polished sections. Laser tracks were positioned to follow along bubble walls in pumice resulting in  $5$  to  $8 \mu m$ -wide troughs with irregular geometries and lengths of  $200 - 500 \mu m$ . For both pyroxenes and glasses, acquired counts per second data were reduced offline and internally normalized to the electron microprobe determined  $SiO<sub>2</sub>$  contents and laser determined <sup>29</sup>Si counts per second. Two-point calibrations were constructed using BCR-2g and either NIST 610 or NIST 612 as standards.

Glass shard separates from the RT were too small for mounting and polishing and their trace element contents were therefore determined using a small-scale dissolution technique on small aliquots (3 to 11 mg) repeated three times for each sample and averaged. Glass separates were dissolved in a two-step digestion process. First, 0.05 ml pure H<sub>2</sub>O, 0.1 ml 37% HCl, 0.1 ml 70% HNO<sub>3</sub>, 0.25 ml 50% HF, and 0.1 ml 71%  $HClO<sub>4</sub>$  were added to each sample vial. Vials were capped overnight and held on a hotplate at 130°C, then uncapped and allowed to evaporate to dryness. Next, 0.05 ml pure H<sub>2</sub>O and 0.1 ml 70% HClO<sub>4</sub> were added, and evaporated to dryness again. The final solutions were gravimetrically diluted to 30 mg in  $2\%$  HNO<sub>3</sub> with 4 drops  $30\%$  H<sub>2</sub>O<sub>2</sub> and 1 drop  $50\%$  HF. Solutions were analyzed on an Element2 mass spectrometer coupled to a Cetac Aridus desolvating nebulizer to increase sensitivity and to reduce oxides. International standards BCR-2, BHVO-1, and in-house standard GMP were used for calibration. A comparison of BCR-2g and an in-house pyroxene standard (TFPX4) analyzed by both laser ablation and solution methods is shown in Fig. 3; there is no significant difference in REE contents analyzed by the two methods.

Gadolinium is not used in this study; there are unresolved issues with the <sup>157</sup>Gd isotope in many of the laser ablation analyses of both pyroxenes and glasses. Scandium contents in BT glasses are below the detection limits of the laser ablation procedure, while doubly charged zirconium  $(^{90}Zr^2)$  in the very Zr-rich FM phonolite glasses interferes with <sup>45</sup>Sc; consequently,  $D_{S_c}$  is only reported for the RT system.

Major elements in FM and RT pyroxenes were measured using a Cameca electron microprobe in the Washington State University GeoAnalytical Laboratory. The instrument was operated at 15kv, using a spot diameter of ~4 microns and a beam current of 12nA.

### **Pyroxene chemistry**

The Fasnia pyroxenes are sodic salites  $(Wo_{49}En_{31}Fs_{20}$  to  $Wo_{45}En_{28}Fs_{27})$ , the BT pyroxenes are ferrohedenbergites to ferroaugites ( $Wo_{44}En_{10}Fs_{46}$  to  $Wo_{40}En_{10}Fs_{50}$ ), and RT pyroxenes are ferroaugites to ferrohedenbergites ( $Wo_{43}En_{20}Fs_{37}$  to  $Wo_{41}En_{6}Fs_{53}$ ) in the pyroxene quadrilateral of Poldervaart and Hess (1951); see Fig. 2. The RT and BT pyroxenes have similar major and trace element contents, but are significantly different from the Na-salites (Table 2 and Figs. 4), which have higher Na, Mg, Al, and Ti contents, and lower Fe, Mn and Si contents.

Compositional differences between hedenbergites and salites result in contrasting major element cation site assignments. Silicon occupies the T-site and Al makes up any Si shortfall in tetrahedral coordination. Si+Al is approximately equal to 2 pfu for all pyroxenes. M1 hosts all  $Ti^{4+}$  and also hosts Fe, Mn, and Mg, but significant Ca+Na shortfalls  $\leq 1$  pfu) in the hedenbergites require that at least some Fe and Mn, and probably Mg, occupy M2. In the Na-salites this substitution is minor.

For the sake of simplicity and consistency, we calculate  $Fe<sup>3+</sup>$  contents for all the pyroxenes using the stoichiometric assumption of Droop (1987). Approximately 50% of total Fe in the FM pyroxenes, 10% of total Fe in the RT pyroxenes, and  $\leq 10\%$  of total Fe in the BT pyroxenes is found to be Fe<sup>3+</sup>. We assume no preferential ordering of Mn, Fe, and Mg into M2 and M1 structural sites in the pyroxenes, rather we calculate sizes for the 6-fold coordinated site from the weighted average of Mn,  $Fe^{2+}$ ,  $Fe^{3+}$ , and Mg contents. All  $Fe<sup>3+</sup>$  is considered to be in 6-fold coordination.

### **Trace element partitioning between pyroxenes and glasses**

#### *Mineral chemistry effects*

The lattice structure of clinopyroxene changes in response to changing Fe/Mg and Na/Ca due to the size differences between Na and Ca, and Fe and Mg (Cameron and Papike, 1980). Since Mn is slightly larger than  $Fe<sup>2+</sup>$  (Shannon, 1976), incorporation of Mn must also result in structural changes. These structural changes should have consequences for trace element partitioning behavior, which should therefore differ between Fe-rich and the more common Mg-rich pyroxenes.

Pyroxene has the general formula  $[M2][M1][T_2O_6]$ . The tetrahedral site (T) is filled with Si and lesser amounts of <sup>IV</sup>Al. The M2 polyhedron is larger and always more distorted than M1 (Cameron and Papike, 1980) and will accommodate VI-, VII- and VIII-coordinated univalent, divalent and trivalent cations with ionic radii from  $\sim$ 0.72 Å (<sup>VI</sup>Mg) to  $\sim$ 1.16 Å (<sup>VIII</sup>Na). Sevenfold coordination has only been described for some pigeonites (Cameron and Papike, 1980) and is not considered further here. A preference for either M1 or M2 is influenced by anion configuration around each site, and by the ionic radius and electronegativity of the cations (Cameron and Papike, 1980). M1 is the smaller of the two cation sites and accommodates divalent, trivalent, tetravalent, and pentavalent cations in 6-fold coordination, with ionic radii between  $\sim$ 0.53 Å (<sup>VI</sup>Al<sup>3+</sup>) and  $\sim$ 0.83 Å (<sup>VI</sup>Mn<sup>2+</sup>). Larger cations such as <sup>VIII</sup>Na (1.18 Å) and <sup>VIII</sup>Ca (1.12 Å) will prefer M2 to M1 (Robinson, 1980).

#### *REE partitioning model*

In the lattice strain model (LSM) of Brice (1975) and Blundy and Wood (1994), the partition coefficient *D*<sup>i</sup> at equilibrium of a cation with radius  $r_i$  entering a particular crystal lattice site M is given by

$$
D_i = D_0^* \exp\{-4\pi E_M N_A [r_0(r_i - r_0)^2/2 + (r_i - r_0)^3/3]/RT\}
$$
 (1)

where  $D_0$  is the partition coefficient of an isovalent ion of radius  $r_0$  which enters the site without strain,  $E_M$ is the Young's modulus of site M,  $N_A$  is Avogadro's Number,  $R$  is the universal gas constant, and  $T$  is temperature in Kelvin. The relationship is most often plotted on an Onuma diagram as the logarithm of the partition coefficient versus ionic radius, producing a parabolic distribution (Fig. 5a). The partition coefficient  $D_0$  for a cation with optimal radius  $r_0$  fixes the apex of the parabola, and its curvature is controlled by the Young's modulus  $(E_M)$  for the site. Larger  $E_M$  values 'tighten' the parabola, reflecting greater stiffness of the site and therefore a relatively narrower range of radii for cations that significantly substitute into the site. For the application of the model below, it is not necessary to know  $r_0$  and  $D_0$  *a priori* as long as there are enough elements to plot to define the shape of the parabola.

Numerous studies have successfully applied the LSM to mineral-melt partitioning, especially for Ca-Mg pyroxenes (Wood and Blundy, 1997; Vannucci et al., 1998; Hill et al., 2000; Wood and Trigila, 2001; Tiepolo et al., 2002). These studies report smooth parabolic distributions of log  $D_i$  vs.  $r_i$  for REE in pyroxenes that can be modeled by assuming that all  $REE^{3+}$  are substituting for eight-fold coordinated  $Ca^{2+}$ . However, iron rich pyroxenes found in evolved felsic liquids (e.g. Larsen, 1979; Wörner et al., 1983;

Mahood and Stimac, 1990; Fedele et al., 2009) show REE patterns with elevated partition coefficients from Er to Lu, rather than a simple parabolic form (Figs. 1, 6). Because heavy REE have ionic radii more similar to Fe and Mn in 6-fold coordination than Ca in eight-fold coordination (Table 3), we propose that elevated  $D_{\text{HREE}}$  in Fe-rich pyroxenes reflects enhanced compatibility relative to lighter REE caused by heavy REE occurring in both eight-fold coordination (substituting for Na or Ca in M2) and 6-fold coordination (substituting for Fe and Mn in M1 or M2). If this is the case, then each coordination state ( $\rm{V^IREE^{3+}}$  and  $\lim_{\text{N}}$ REE<sup>3+</sup>) is governed by its own LSM equation, and mineral/melt *D* values are the sum of the two parabolic functions (Figs. 5 & 7).

#### *Identification of pyroxene-glass pairs*

Before these relationships can be explored, it is necessary to establish that the pyroxene grains and glasses used to obtain *D* values represent equilibrium pairs. One possible problem is late crystallization of accessory phases that concentrate REE and would therefore deplete the melt in REE after pyroxene growth. While we cannot completely rule out this effect, we note that it would tend to deplete middle REE in FM glasses and light REE in the rhyolitic glasses, thus enhancing the apparent *D* values for these elements with respect to heavy REE, the opposite of what we in fact observe. In fact, the FM pyroxenes have a deeper middle REE trough than do the glasses (Fig. 3).

In all three examples, pyroxene grains are euhedral and appear to be in textural equilibrium with the enclosing glass. However, in the case of the compositionally variable Fasnia phonolite, magmatic processes have scrambled crystals and glasses (Olin, 2007), precluding confident identification of mineralglass pairs based on physical proximity or grain shape. The method outlined below is independent of major element mineral compositions, assumes that the LSM is a good first-order predictor of trace element partition coefficients, and does not rely on textural relationships. A crucial assumption is that the light and middle REE  $(La - Dy)$  are not partitioned between 6-fold and 8-fold sites, but reside exclusively on the 8fold site. We emphasize that, qualitatively, our finding of elevated pyroxene/melt partition coefficients for heavy REE remains robust even if mineral-melt pairs are randomly assigned within each of the three systems. We make the effort to identify equilibrium pairs in order to isolate and quantify the heavy REE enrichment in the pyroxenes to the maximum extent allowed by the data.

For 25 FM Na-salite and 44 BT hedenbergite grains, apparent  $D_{\text{REE}}$  values were calculated using LA-ICP-MS-determined concentrations of La, Ce, Pr, Nd, Sm, Tb and Dy (approximately 500 phonolite glasses in the case of the Fasnia and approximately 100 high-silica rhyolite glasses from the BT). Eu is omitted due its unconstrained oxidation state, and heavy REE (Ho – Lu) for the reasons outlined above. Equilibrium between the pyroxenes and the different glasses is evaluated by assuming that isovalent cations entering a single lattice site will always follow a parabolic distribution in  $\log D_i$  vs.  $r_i$ , using equation (1) rewritten as:

$$
[RT/(-4\pi N_A)] * [\ln(D_i/D_0)] = E_M[r_0/2^*(r_i - r_0)^2 + 1/3^*(r_i - r_0)^3]
$$
 (2)

If the data follow the parabolic function exactly, the result is a straight line with slope  $E_M$  (Fig. 5c). Using equation (2) with the selected REE<sup>3+</sup> La to Dy, the pyroxene-glass pair with the best  $R^2$  correlation coefficient for the linear regression represents the closest approach to equilibrium and yields pyroxene-melt partition coefficients. We applied a filter of  $R^2 > 0.95$  to 12,500 possible pyroxene-glass pairings from the FM and 6,400 pairs from the BT, yielding 300 and 220 near-equilibrium pairs respectively. This approach is more rigorous than simply using predicted absolute *D* values because the fit does not depend on knowing  $D_0$  or  $r_0$ , and relies on multiple elements. For the 40 RT pyroxenes, partition coefficients were calculated using pyroxenes paired with equilibrium glasses identified by Streck and Grunder (1997) and re-analyzed for this study.

From equation (2), the resultant calculated  $\frac{VIII}{E^3}$ <sub>M2</sub> values for the Na-salites are 113 - 240 GPa (ave. 177) GPa), for BT hedenbergites 306 - 387 GPa (ave. 339 GPa) and for the RT ferroaugites 249 to 358 GPa (ave. 296 GPa). Our estimates of  $\frac{VIII}{E^3}$ <sub>M2</sub> for the BT and RT pyroxenes compare well with published values of 256 to 387 GPa (Hill et al., 2000; Lundstrom et al., 1998). Those for the FM Na-salites are lower due to the influence of Na, a result also found experimentally by Bennett et al. (2004).

Model curves fitted to the average pyroxene-glass partition coefficients for each of the three populations are plotted on Figure 7. Parameters for each model curve are given in Table 4. Temperatures are from the literature for each system: Fasnia, assumed similar to other Tenerife phonolites, 825°C (Andujar et al., 2008); RT, 800°C (Streck and Grunder, 1997); and BT at 700°C (Warshaw and Smith, 1988). Values for the other parameters used in the model curves are determined in a two-step approach. First, values for the eight-fold coordination LSM equation come from equilibrium pyroxene-glass pairs identified using LREE and MREE in equation (2). This results in eight-fold coordination model values of:  $VIII_{r_0} = 1.06$  Å,  $VIII_{D_0} =$ 1.6, and <sup>VIII</sup> $E^{3+}_{\text{M2}} = 177$  GPa for Na-salite; <sup>VIII</sup> $r_0 = 1.055$  Å, <sup>VIII</sup> $D_0 = 5.2$ , and <sup>VIII</sup> $E^{3+}_{\text{M2}} = 296$  GPa for RT hedenbergites; and <sup>VIII</sup> $r_0 = 1.055$  Å, <sup>VIII</sup> $D_0 = 9$ , and <sup>VIII</sup> $E^{3+}_{M2} = 339$  GPa for BT hedenbergites. Then, plausible values are chosen for the six-fold coordination LSM equation to fit the cumulative (VIII-fold + VI-fold) curves through the data, i.e., the heavy REE. The value of  $V^I_{r0}$  = 0.724 Å, 0.758 Å and 0.772 Å for FM, RT and BT pyroxenes, respectively, is derived from the combined effective ionic radius, or weighted average, of the  $Fe^{2+}$ ,  $Fe^{3+}$ , Mn, and Mg contents, and allows the model to be reasonably fitted through the  $D_{\text{HREE}}$  and the notably high  $D_{\text{Sc}}$  values for RT hedenbergites. For FM Na-salites  $^{VI}D_0 = 500$  and  $^{VI}\overline{E}^{3+}_{\text{M}} =$ 775 GPa, for RT hedenbergites  $^{VI}D_0 = 150$  and  $^{VI}E^{3+}_{M} = 1000$  GPa, and for BT hedenbergites  $^{VI}D_0 = 200$ and  ${}^{VI}E^{3+}{}_{M}$  = 1100 GPa. The  ${}^{VI}E^{3+}{}_{M}$  values found here are comparable to values determined for diopside by Hill et al. (2000), Lundstrom et al. (1998) and Hazen & Finger (1979). It should be noted that <sup>VI</sup>D<sub>0</sub> values are quite sensitive to small changes in  $V_{T_0}$  in the sixfold model equation. This is a function of the higher  $V_{T_{\text{EM}}}$  values for this coordination state (Hill et al., 2000).

Immediately apparent from this exercise is that the negative  $D<sub>Y</sub>$  anomaly seen in FM, RT, BT, and other iron-rich pyroxenes (Figs. 1,7) is not explained by the LSM, which predicts a smooth pattern across Dy, Y and Ho. This feature is discussed below.

#### *Trace element partition coefficients*

Each of the three Fe-rich clinopyroxene populations studied here exhibits a significant range in  $D_{REF}$ , (Figs.  $6 - 9$ ; Table 5). Na-salites have the following ranges in D-values for selected REE and Y (YREE):  $D_{La}$  0.4 to 1.1 (aver. 0.6),  $D_{Sm}$  1.0 to 2.4 (aver. 1.5),  $D_{Lu}$  3.5 to 5.9 (aver. 4.5), and  $D_{Y}$  0.9 to 1.8 (aver. 1.3). RT hedenbergites have  $D_{La}$  0.7 to 1.2 (aver. 1.0),  $D_{Sm}$  2.9 to 6.6 (aver. 5.0),  $D_{Lu}$  4.0 to 7.9 (aver. 6.0), and  $D_Y$  2.1 to 4.8 (aver. 3.7). BT hedenbergites have  $D_{La}$  0.9 to 2.3 (aver. 1.2),  $D_{Sm}$  6.3 to 15 (aver. 8.0),  $D_{Lu}$  7.1 to 14 (aver. 10), and  $D_y$  3.8 to 7.2 (aver. 4.8).  $D_{REE}$  are positively correlated with each other as indicated by the uniform increase of D-values in Figure 5. In general, and especially in ferrohedenbergites from the BT, Y is less compatible into pyroxene than heavy REE of similar radius (Figs. 7, 8); this behavior is discussed below. Calculated D values for LREE and MREE are generally comparable to published values from a range of magma compositions (e.g., Sisson, 1991; Hart and Dunn, 1993; Schosnig and Hoffer, 1998; Hill et al., 2000; Bennett et al., 2004; Huang et al., 2006; Lofgren et al., 2006), but values for heavy REE are significantly higher relative to middle REE. Comparison of FM and RT pyroxenes shows that  $D_{L}$  increases with increasing  $Fe^{2+}$  and Mn contents. Among the RT hedenbergites,  $D_{Lu}$  is however negatively correlated with Fe contents, but positively correlated with Mn contents (Fig. 9).

Partition coefficients for HFSE are shown in Figure 8 and Table 5. Na-salites have the following selected  $D_{HFSE}$ :  $D_{Ti}$  0.8 to 1.9 (aver. 1.3),  $D_{Zr}$  0.7 to 1.7 (aver. 1.2), and  $D_{HF}$  1.0 to 2.8 (aver. 1.8). RT hedenbergites have  $D_{Ti}$  1.2 to 2.0 (aver. 1.5),  $D_{Zr}$  0.16 to 0.26 (aver. 0.21), and  $D_{HF}$  0.22 to 0.37 (aver. 0.28). BT hedenbergites have  $D_{Ti}$  1.3 to 2.0 (aver. 1.6),  $D_{Zr}$  0.17 to 0.40 (aver. 0.26), and  $D_{HF}$  0.27 to 0.98 (aver. 0.44). In all pyroxenes studied  $D_{H_1} > D_{Zr}$  and  $D_{H_0} > D_Y$ , and in the Na-salites  $D_{Ta} > D_{Nb}$  (Ta contents in BT and RT pyroxenes were below detection limits). Calculated HFSE partition coefficients are generally comparable to other published values (e.g., Sisson, 1991; Hart and Dunn, 1993; Schosnig and Hoffer, 1998; Hill et al., 2000; Bennett et al., 2004; Marks et al., 2004; Huang et al., 2006; Lofgren et al., 2006).

#### **Discussion**

#### *Factors controlling HREE enrichment in pyroxene*

The primary factors controlling the style of trace element partitioning between clinopyroxene and silicate melt are mineral composition and the ionic radius of the substituting cation. The generally lower  $V^{III}E^{3+}$ values for Na-salites relative to the ferrohedenbergites and ferroaugites is attributed to their higher Na contents ; Na cations have the effect of opening the calcic pyroxene structure and lowering  $V^{III}E^{3+}$  (Bennett et al., 2004).

We have shown that the HREE enrichment of these pyroxenes is plausibly explained by their high Fe contents (i.e.,  $Fe^{2+}$ ) that are common to all three groups studied here, but variations in  $D_{HREF}$  within the FM and RT pyroxene populations are not apparently controlled by Fe contents alone (Fig. 9); considering the RT pyroxenes in isolation, *D*Lu is strongly positively correlated with Mn contents, but negatively correlated with Fe contents. This suggests that Mn, in fact, has a more powerful effect than  $Fe^{2+}$  in creating the space for 6-fold coordinated heavy REE to substitute into the pyroxene structure. This is consistent with Mn having the larger ionic radius.

#### *Fractionation of Y/Ho, Zr/Hf, and Nb/Ta*

The elements in each of these pairs have closely similar radii and have been viewed in the past as chemically identical at magmatic temperatures. Thus, Y was commonly used as a proxy for the HREE between Tb and Yb prior to the routine availability of high-quality ICPMS data for all 14 REE (for example, Thompson et al., 1984). Significant fractionation in these pairs of elements is observed at submagmatic temperatures (Bau, 1996, 1999; Bau and Dulski, 1999) and, in plutonic rocks, has been ascribed to the action of aqueous fluids during the final stages of crystallization (Irber, 1999). Carbonatite liquids have also been implicated in Zr/Hf and Nb/Ta fractionation (Rudnik et al., 1993). Nonetheless, it is now well established that these element pairs are fractionated during common crystallization-differentiation processes. Magmatic fractionation of Nb from Ta was first demonstrated by Wolff (1984) in Tenerife phonolites and shown to be due to titanite crystallization; Ti-rich phases in general have high *D* values for HFSE and are effective in separating these two elements (Green and Pearson, 1986). The restricted *P-T-X* conditions under which phases such as titanite and rutile are stable has lead to considerable attention being paid to variations in Nb/Ta, in particular, among different terrestrial reservoirs (Rudnik et al., 1993; Collerson et al., 2007; Horng and Hess, 2000; Wade and Wood, 2001; Foley et al., 2002). With the advent of high-quality mineral-melt partitioning data, it has become clear that common phases such as clinopyroxene and amphibole also fractionate Y from Ho, Zr from Hf, and Nb from Ta (Lundstrom et al., 1998; Tiepolo et al., 2002).

The ability of the LSM to predict these fractionations depends on accurate ionic radius data. Nb<sup>5+</sup> and  $Ta<sup>5+</sup>$ are conventionally viewed as having identical radii,  $Zr^{4+}$  and  $Hf^{4+}$  are slightly different, while the radius of  $Y^{3+}$  lies between those of Dy<sup>3+</sup> and Ho<sup>3+</sup> in eight-fold coordination, but is nearly identical to Ho<sup>3+</sup> in 6-fold coordination (Shannon 1976; relevant values are reproduced in Table 3. Note that in the older compilation of Whittaker & Muntus, 1970, the ionic radii of  $Y^{3+}$  and  $Ho^{3+}$  are equal in both 6-fold and 8-fold coordination). Accuracy of the LSM is difficult to assess for Nb-Ta and Zr-Hf, because the very high *E* values associated with high valencies and the small M1 site in clinopyroxene imply easily detectable changes in *D* at variations in  $r \le 0.01$  Å, beyond the resolution of the Shannon (1976) ionic radius data. Blundy and Wood (2003) in fact propose that the ionic radius of  $Ta^{5+}$  is 0.01 – 0.02 Å smaller that that of  $Nb<sup>5+</sup>$ , based on  $D_{Nb}/D_{Ta}$  variations in Ti-pargasite and kaersutite documented by Tiepolo et al. (2002). However, inspection of Figure 7 shows that  $Y^{3+}$  cpx/melt partitioning deviates from that of the HREE to an extent that the LSM cannot explain unless the Shannon (1976) value of the ionic radius for  $Y^{3+}$  is grossly in error, and really lies within the range of ionic radii of the light REE! Titanite displays similar behavior to a greater degree (Olin, 2007). This result is not entirely unexpected; there can be little doubt that ionic radius and charge are more important than other factors in predicting D values, and so the effects of deviations from the LSM should be most apparent among isovalent cations of similar radii. Although small differences in ionic radius may indeed cause fractionation of Nb from Ta and Zr from Hf, this cannot explain the fractionation of Y from the heavy REE. The pairs Y-Ho, Zr-Hf and Nb-Ta share a common feature, which is that the lighter cation has the outer electronic configuration [Kr], and the heavier cation has buried 4*f* electrons and the outer electronic configuration [Xe]. Contrasts in behavior between the members of each pair may therefore be due to electronic effects, such as the related properties of electronegativity and ionic polarizability, and/or mass effects.

#### *Sources of deviation from the LSM*

The LSM assumes that cations behave as 'billiard balls' with evenly distributed surface charge. Deviations are expected where outer electronic structure cannot be so modeled, for example crystal field effects resulting from the involvement of *d* orbitals in bonding, or the presence of a lone pair of electrons (Blundy and Wood, 1994). The ions  $Y^{3+}$ ,  $Ho^{3+}$ ,  $Zr^{4+}$ ,  $Hf^{4+}$ ,  $Nb^{5+}$  and  $Ta^{5+}$  all have inert gas outer electronic configurations and therefore approximate to charged billiard balls, but exhibit variation in electronegativity and ionic polarizability (Table 6), i.e. the tendency to develop uneven charge distribution.

Horng and Hess (2002) invoked the polarizability difference of  $Nb_2O_5$  and Ta<sub>2</sub>O<sub>5</sub> (as a proxy for NbO<sub>6</sub> and TaO<sub>6</sub> octahedra in minerals and melts), calculated from specific refractivity values tabulated by Jaffe (1988), to explain experimentally-determined fractionation of Nb/Ta between rutile and haplogranite melt. Using the same tabulation (Table 6), polarization effects are predicted to operate in one sense for Y-HREE and Zr-Hf, and in the opposite sense for Nb-Ta. Using the more recent compilation of ionic polarizabilities of Shannon and Fischer (2006), dipole effects for Y-HREE should act in the opposite sense to Zr-Hf and Nb-Ta (Table 6). The differences in electronegativity and polarizability between Y and HREE are also small. Although not consistent in the prediction of Zr-Hf behavior, both approaches predict opposite senses of Y-Ho and Nb-Ta partitioning. This contrasts with our data, which shows that the heavy element in each pair is always the more compatible. With the understanding that as yet undetected differences in ionic radius may trump secondary effects, we briefly speculate on the possible role of mass effects.

### *Mass fractionation effects?*

By analogy with the theoretical treatment of stable isotope fractionation (Chacko et al., 2001), mass fractionation between two otherwise identical ions may be crudely approximated by treating the metaloxide bond as a diatomic molecule that behaves as a simple harmonic oscillator (spring) connecting two masses  $m<sub>M</sub>$  and  $m<sub>O</sub>$  where M and O denote the metal and oxygen respectively. The vibrational frequency,  $\nu$ , of the oscillator is given by

$$
v = (k/\mu)^{0.5}/2\pi \tag{3}
$$

**halo also solute Equal to the spring (a property related to the elastic modulus), and**  $\mu$  **is the** reduced mass, given by  $\mu = m_M m_0/(m_M + m_0)$ . In this model, a strong chemical bond can be thought of as a spring with a high force constant, resulting in a higher vibrational frequency. From equation (3), changing the mass of the metal ion gives

$$
v^*/v = \left(\mu/\mu^*\right)^{0.5} \tag{4}
$$

where the asterisk denotes substitution of the heavier atom. Due to the inverse relation between frequency and mass, the heavy molecule has a lower vibrational frequency than the lighter molecule and hence is more energetically stable. Partitioning between two phases then depends on their competition for the heavy atom. Equation (4) can be rearranged to yield

$$
v - v^* = v[1 - (\mu/\mu^*)^{0.5}] \tag{5}
$$

from which it can be seen that the heavy atom favors the phase with stronger M-O bonds (higher  $v$ ).

Because we have no information on the strengths of trace metal – oxygen bonds in silicate liquids, we cannot make quantitative estimates of pyroxene-melt mass fractionation effects for trace elements. Nonetheless, some qualitative predictions are possible based on equations  $(3) - (5)$ . First, the sign of

heavy/light partitioning should always be the same for any given phase/melt pair, i.e., either the heavy or the light element should prefer the solid phase for each of the pairs Y-Ho, Zr-Hf, and Nb-Ta. Second, we expect that the mineral phase will in most cases favor the heavy element in each pair. This is because M-O bonds in silicate liquids are assumed to be considerably weaker than the same bonds in solid phases; the very success of the LSM, which ignores melt structural effects, in predicting first-order partitioning behavior supports this assumption. Our data, plus a review of the literature for pyroxene-melt, amphibolemelt, titanite-melt, rutile-melt, magnetite-melt, and ilmenite-melt partitioning of Y/Ho, Zr/Hf and Nb/Ta reveals that in each case the partition coefficients are higher for the heavier elements (e.g., Stimac and Hickmott, 1994; Bottazzi et al., 1999; Oberti et al., 2000; Horng and Hess, 2000; Nielsen and Beard, 2000; Tiepolo et al., 2002; Prowatke and Klemme, 2005). This would be a remarkable coincidence if small differences in radii were solely responsible, given the variety of phases and lattice sites represented. Third, the discrepancy between partitioning behavior of the heavy and light element in each pair should increase with ionic charge. This is because  $(n - n^*)/n$  values (eqn. 5) are approximately constant (0.0365  $\pm$  0.007) for each of the three element pairs, but M-O bond strength is expected to increase with increasing charge on M, all else equal, and so the mass fractionation effect should be enhanced. This effect is seen in our data (Fig. 8). Finally, mass fractionation effects are expected to be inversely correlated with temperature. This effect is also seen in our data (Figure 8), for which  $T_{\text{Fasnia}} > T_{\text{RT}} > T_{\text{BT}}$ .

Niu & Hékinian (1997) and Niu (2004) also appealed to mass-dependent fractionation as a possible explanation for large variations in Zr/Hf and Nb/Ta among abyssal peridotites, and noted that kinetic effects enhance mass fractionation because diffusion coefficients are themselves mass dependent. It is possible that, at the relatively low magmatic temperatures of the systems investigated here, chemical kinetics could play a detectable role in cation partitioning between pyroxene and melt, and may be more marked than equilibrium mass fractionation effects. Variable Y/Ho, Zr/Hf and Nb/Ta remains an intriguing problem in igneous geochemistry that awaits a fundamental explanation.

#### **Conclusions**

Heavy REE enrichment in Fe-rich clinopyroxenes from felsic magmas is caused by substitution into 6 coordinated sites (M2 or M1). Increased heavy REE compatibility coincides with the higher Fe and Mn contents of these pyroxenes relative to more magnesian pyroxenes described elsewhere. This behavior can be modeled by summing the results for the two lattice strain model equations for 6-fold and 8-fold coordinated 3+ cations with reasonable parameter values. Fractionation of Y from REE is observed in all three systems studied here and is not predicted by the lattice strain model. This feature cannot be explained by cation size differences, and may instead be a consequence of electronic effects or, as we speculate here, mass-dependent fractionation. If so, the same effect(s) may possibly play a role in the fractionation of Zr from Hf, and of Nb from Ta.

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# **Figure captions**

**Figure 1:** Clinopyroxene-melt REE partition coefficients versus atomic number of pyroxenes from the literature and from this study. Literature data are from pyroxenes in alkali basalt (Hart and Dunn, 1993; Mg# 55.2; black triangles), lherzolite (Blundy et al., 1998; Mg# 83.7; black diamonds), phonolite (Wörner et al., 1983; Mg# 65.6; black circles), phonolite (Larsen, 1979; Mg# 25.3; black squares), pantellerite (Mahood and Stimac, 1990; Mg# 18.3; open triangles), trachyphonolite (Fedele et al., 2009; Mg# 67; open circles), high-silica rhyolite (Sisson, 1991; Mg# 16.8; open squares), rhyolite (Nash and Crecraft, 1985; Mg# 41.7; dashed x). Data from this study; average Na-salite in Fasnia Member phonolite (Mg# 48 to 59; red), average ferroaugite/ferrohedenbergite in Rattlesnake Tuff rhyolite (Mg# 9.6 to 32; solid black), average ferroaugite/ferrohedenbergite in Bandelier Tuff high-silica rhyolite (Mg# 14.2 to 16.2; blue).

**Figure 2:** Pyroxene quadrilateral showing Tenerife Na-salites (FM), Rattlesnake Tuff (RT) ferrohedenbergites to ferroaugites, and Bandelier Tuff (BT) ferrohedenbergites. Pyroxene quadrilateral after Poldervaart and Hess, 1951.

**Figure 3:** Rare earth elements contents of average pyroxene (solid symbols) and glass (open symbols) compositions for FM (red), RT (black), and BT (blue) normalized to Bulk Silicate Earth, McDonough & Sun, 1995. For comparison, values determined for standard materials BCR (gray) and in-house pyroxene (TFPX, black) are also shown for laser ablation (solid, no symbols) and solution (dashed, no symbols) methods.

**Figure 4:** Covariations of selected minor and trace elements and ratios in the pyroxenes.

**Figure 5:** Lattice strain model plotted (a) as Log(partition coefficient) versus ionic radius for idealized elements of the same valence state and same coordination number, where  $D_0$ is the partition coefficient for the optimal ionic radius  $r_0$ , and  $E_M$  is the Young's modulus for the site, (b) showing  $REE<sup>3+</sup>$  partition coefficient distribution for six-fold and eightfold coordination in pyroxene, and (c) shows equation (2), the linear expression of model equation. Samples shown in (c) have the best  $R^2$  values from each of the three populations (FM red, RT black, BT blue), for elements La through Dy in eight-fold coordination (see text for details).

**Figure 6:** Calculated REE partition coefficient ranges for FM, RT, and BT pyroxenes.

**Figure 7:** Model curves for  $REE<sup>3+</sup>$  cations for FM (a), RT (b), and BT (c) pyroxenes. The model curves are for elements in six-fold (red, squares) and eight-fold (blue, diamonds) coordination and the sum of the two curves (solid black, triangles). The average composition for each population is also shown (black dashed, gray x), but is obscured by the model curves. Note that the negative Y partition coefficient anomaly that is especially prominent in average RT and BT hedenbergites, and in other iron-rich pyroxenes in Figure 1, is not predicted by the model (see text for details).

Figure 8: Partition coefficient co-variation and ratio plots for selected trace elements.

**Figure 9:**  $D_{\text{Lu}}$  versus  $Fe^{2+}$  and Mn contents of the pyroxenes. BT samples are omitted from this plot because trace element analyses do not correspond to locations for major element determinations.



**Figure 1**



**Figure 2**

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**Figure 3**

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**Figure 6**



**Figure 7**

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**Figure 8**

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**Figure 9**

# **Table captions**

**Table 1.** Characteristics of the three felsic volcanic units analyzed in this study.

**Table 2:** Representative major and trace element contents of FM, RT and BT pyroxenes and glasses. BT pyroxene major element and trace element values are from different grains.

**Table 3:** Ionic radii for major and trace element cations in eight- and six-fold coordination, from Shannon (1976).

**Table 4:** Values for lattice strain model parameters used to generate model curves in Figure 7.

**Table 5:** Pyroxene-glass partition coefficients; maximum, minimum, average, and model curve values.

**Table 6:** Electronic polarizability, α, of Y, HREE, and HFSE cations, from Shannon and Fischer (2006).  $\Delta = [\alpha_{\text{(light ion)}} - \alpha_{\text{(heavy ion)}}]/\alpha_{\text{(light ion)}}$ 



\*The BT pyroxenes and glasses in this study are taken from the first  $120 \text{ km}^3$  of homogeneous magma ejected during the Otowi Member eruption. The total volume for the unit is  $\sim$ 450 km<sup>3</sup>.

\*\*NBO/T =  $(2O + 4T)/T$ , where O = total oxygen atoms and T =  $(Si + A1 + Fe^{3+} + Ti)$  on a molar basis, and assuming  $3\%$  dissociated  $H_2O$ , corresponding to total magmatic water contents of 5% in each case (Sowerby & Keppler, 1999).



# **Table 3**



# **Table 4**





# **Table 6**

